

# Quantum gases in Optical Lattices

## Lecture #3

**David CLEMENT**

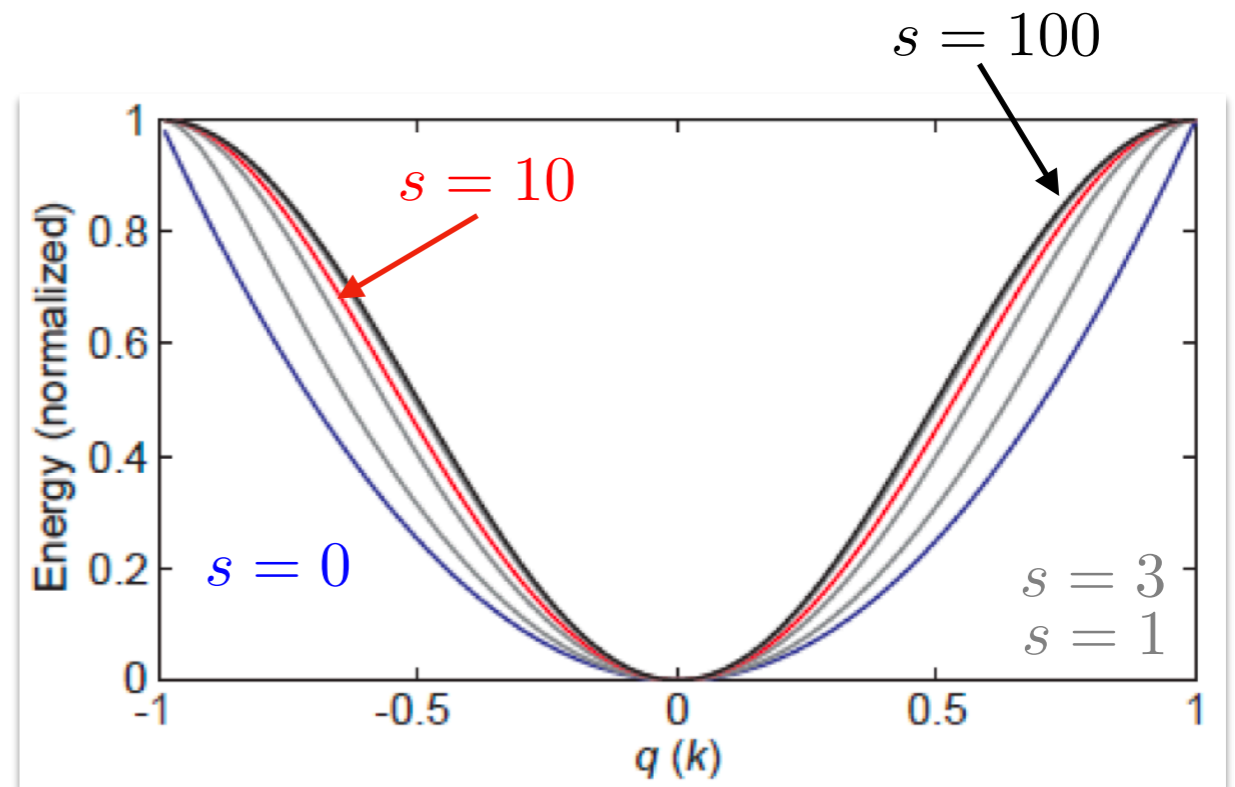
*Institut d'Optique - Palaiseau, France*

# Lowest-energy band

Normalized and shifted lowest-energy bands.

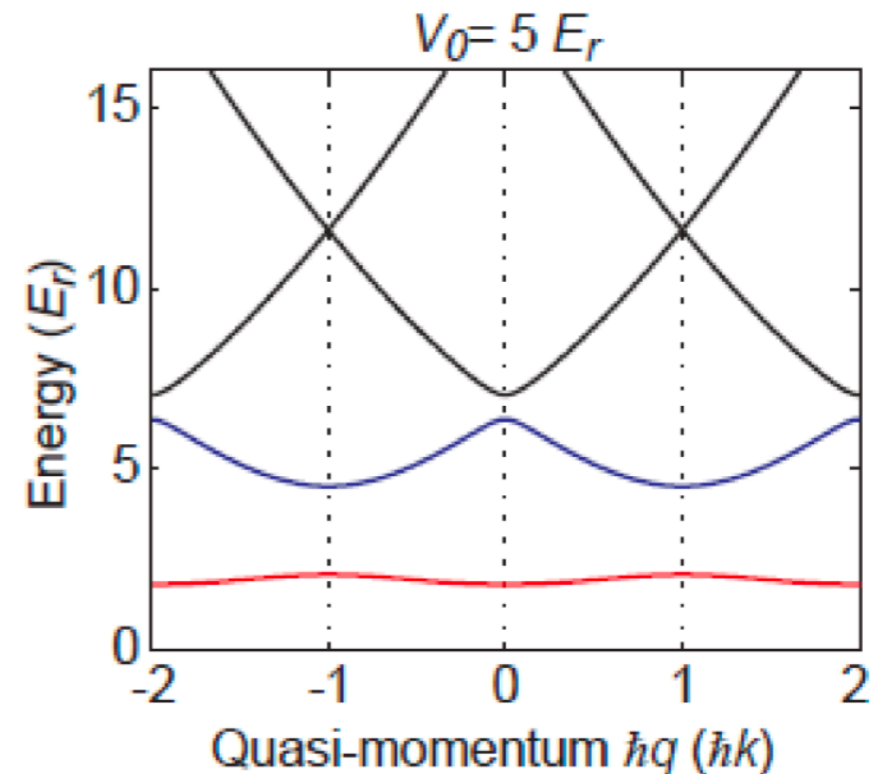
lowest-energy band dispersion smoothly changes from parabolic to sinusoidal:

$$E(q) - E_0 \simeq 4J \sin(qa/2)^2$$

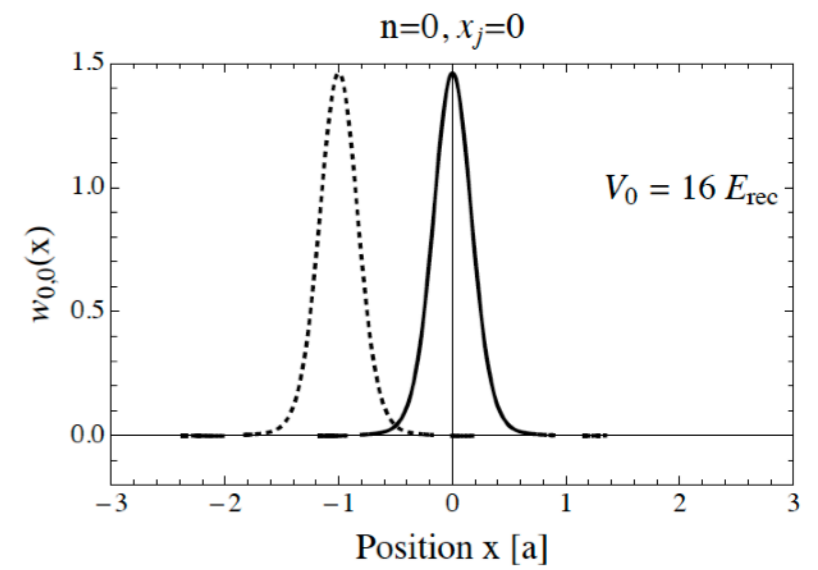
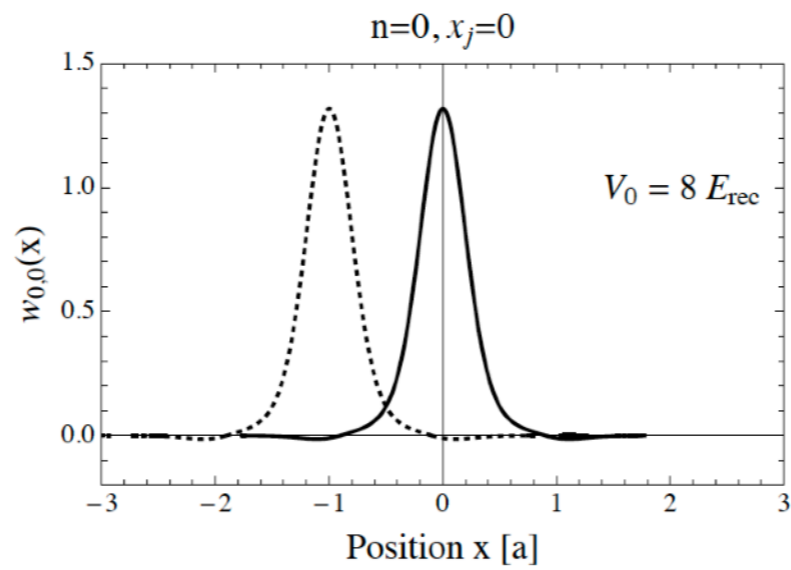
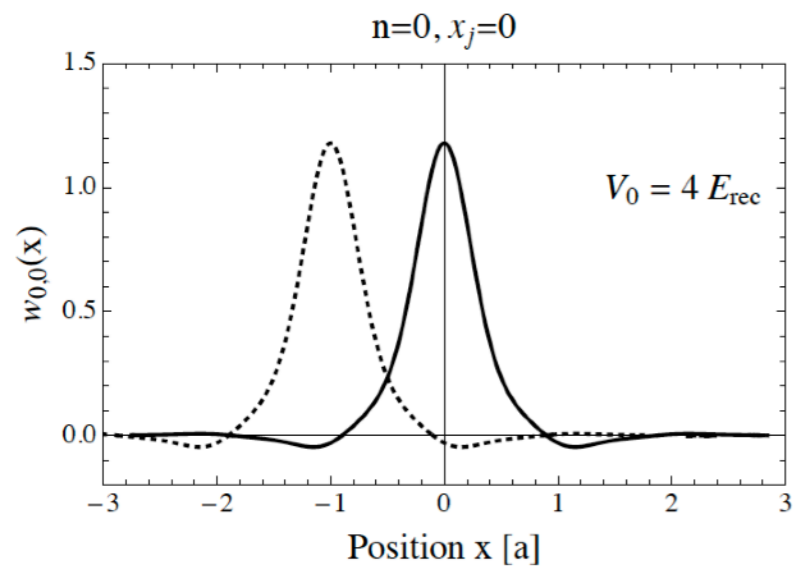
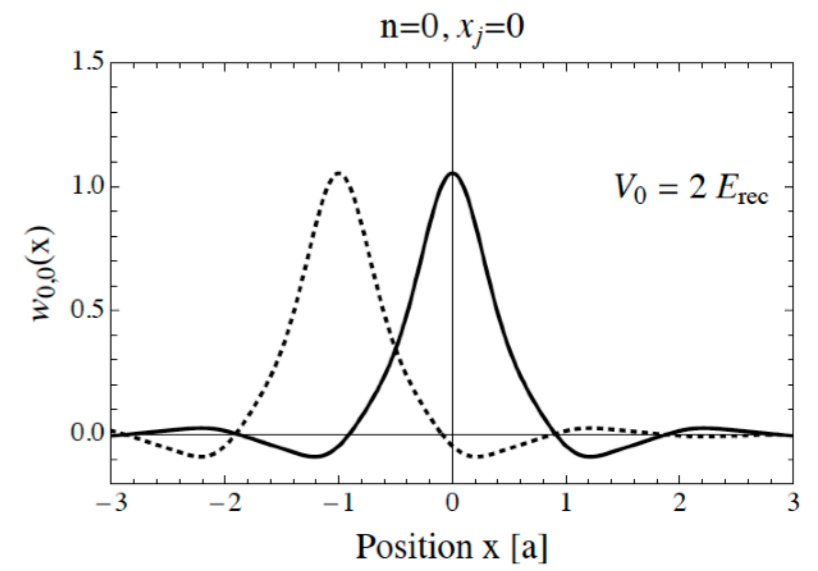
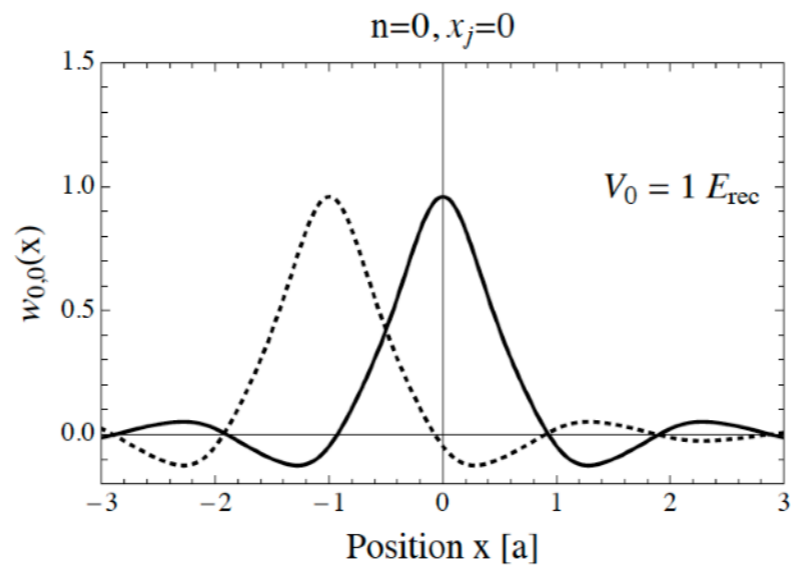
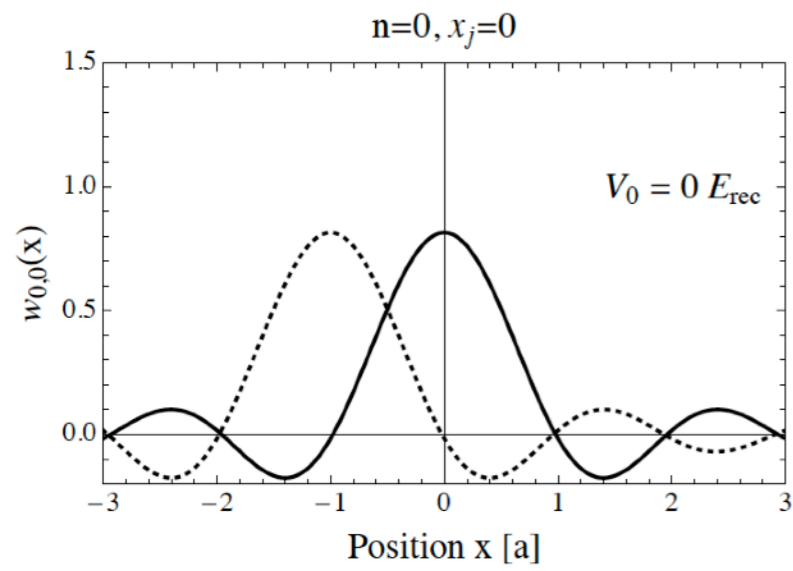


**Tight-binding regime** (*lectures #3, #4*)

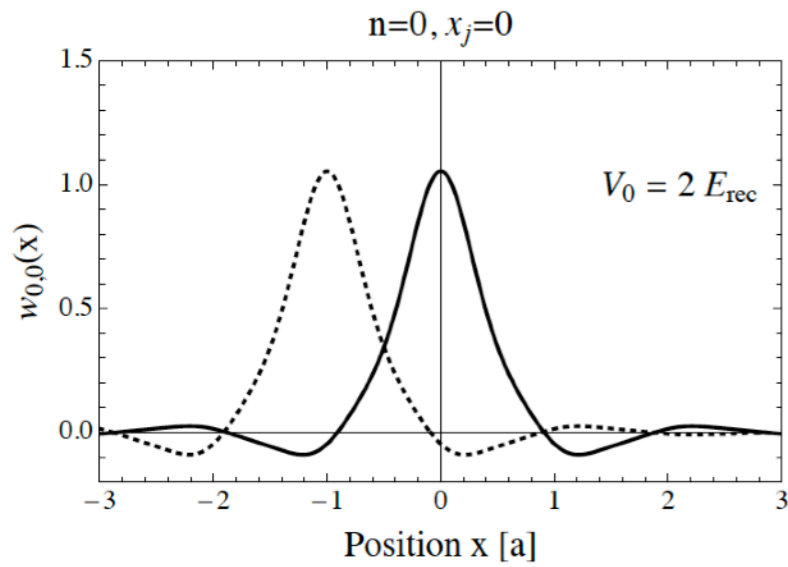
$s \gg 5$  : energy bandwidth  $\ll$  energy bandgap



# Wannier functions



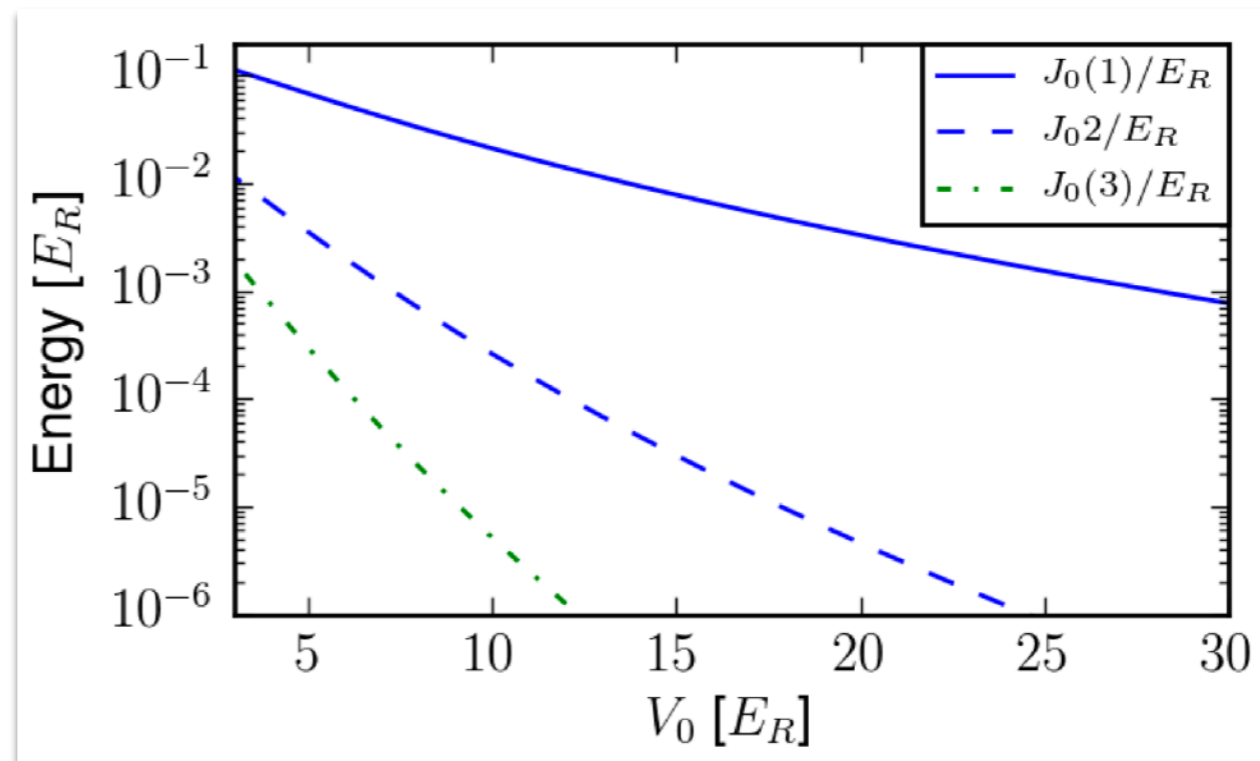
# Tunnelling in the lattice



$$J = J_0 = \int dx \omega_0^*(x+a) \left( -\frac{\hbar^2}{2m} \Delta_x - V_{\text{lat}}(x) \right) \omega_0(x)$$

**Tight-binding regime:**

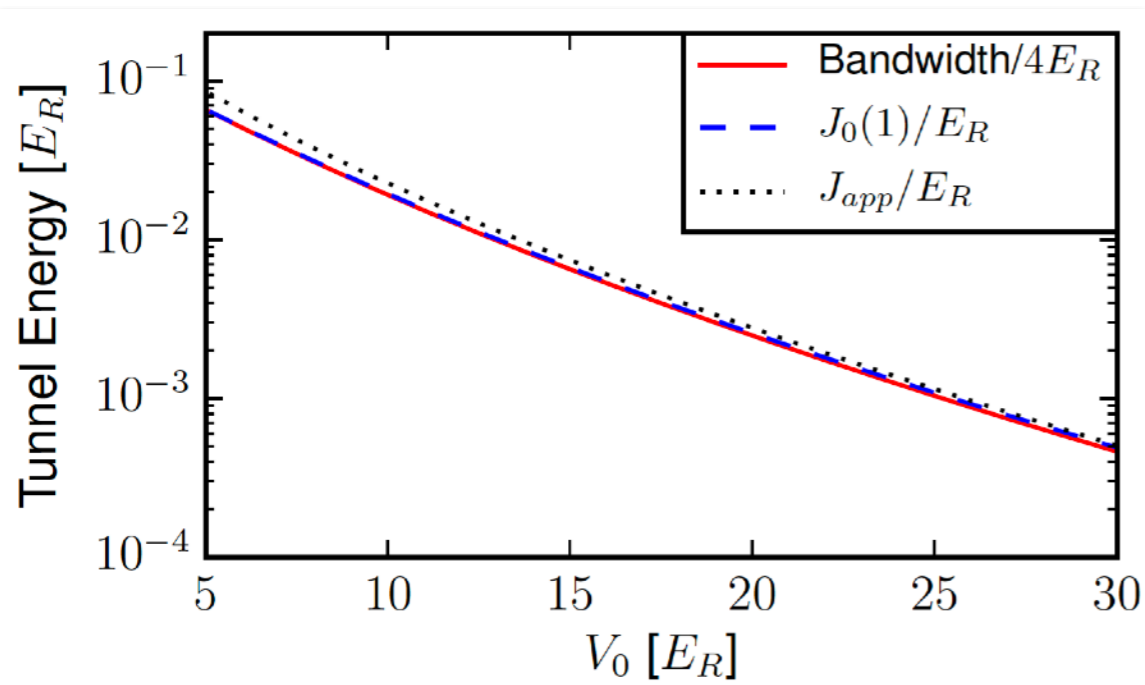
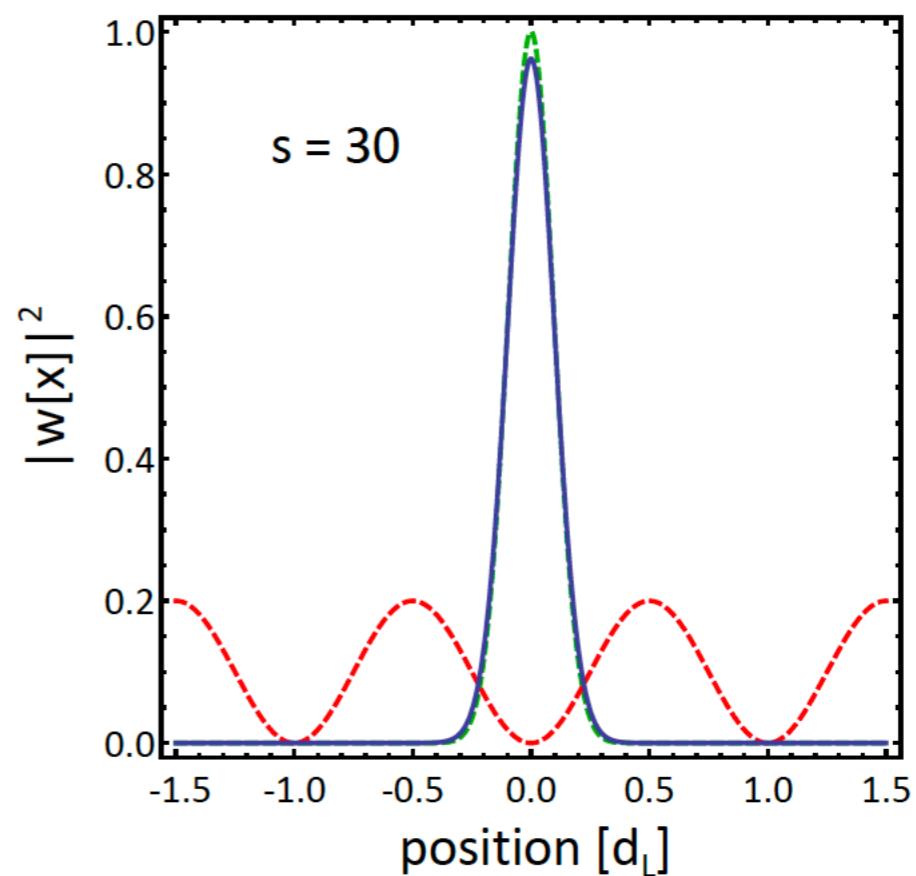
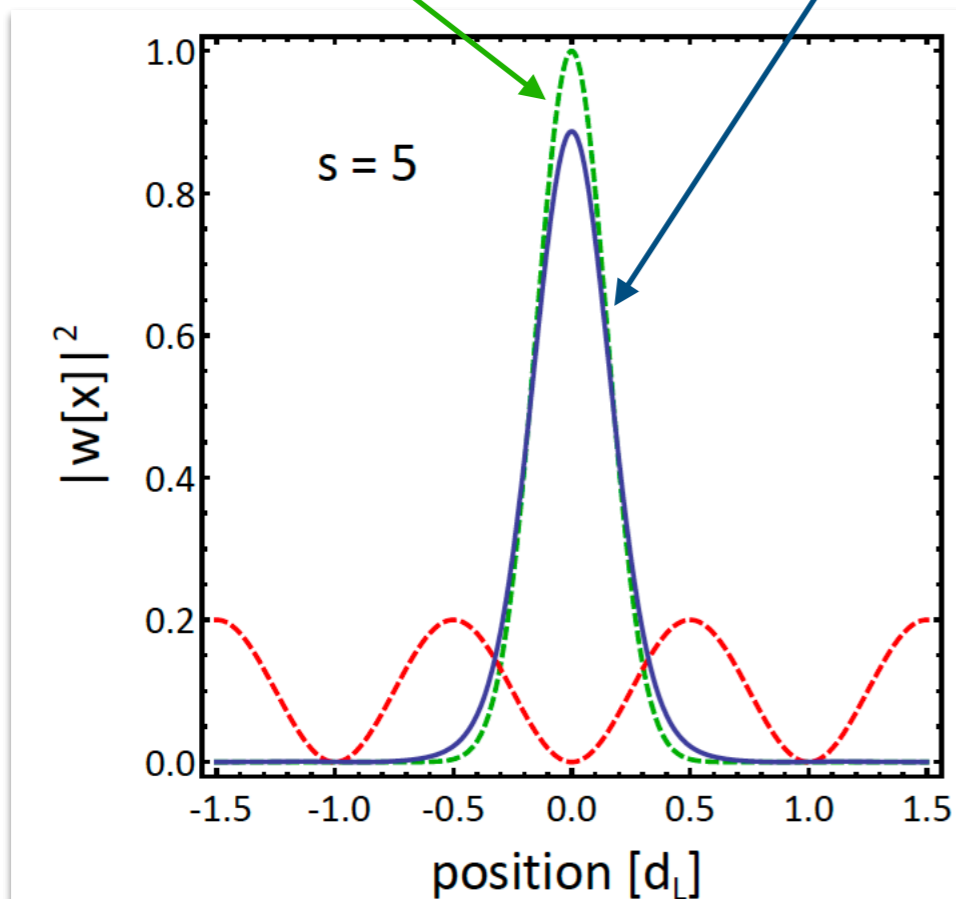
only nearest neighbour tunnelling



# Deep lattices: harmonic approximation

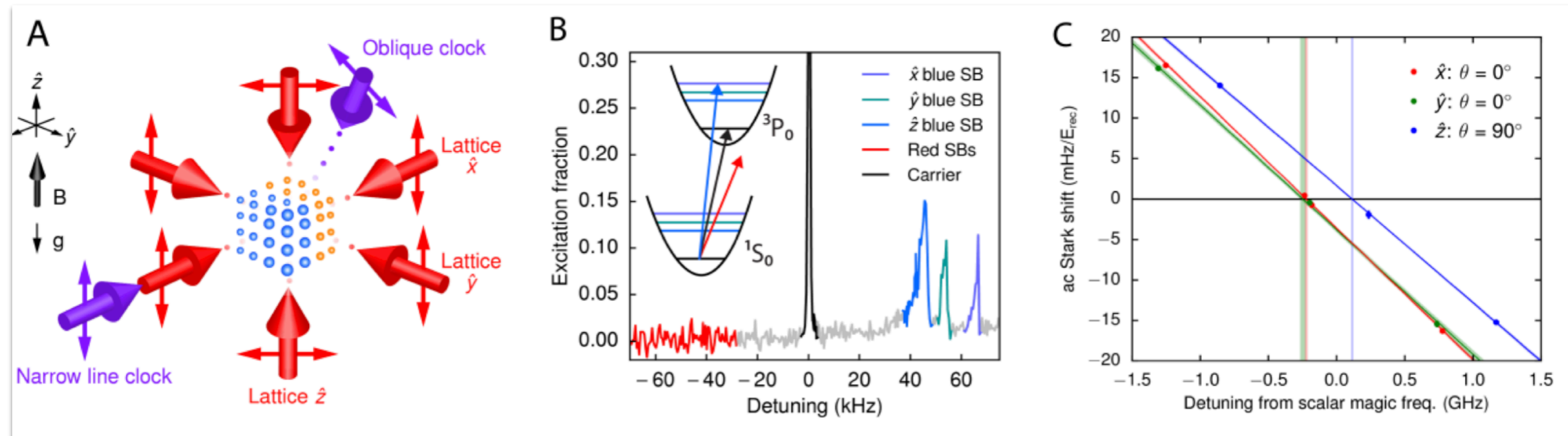
harmonic approximation

Wannier function

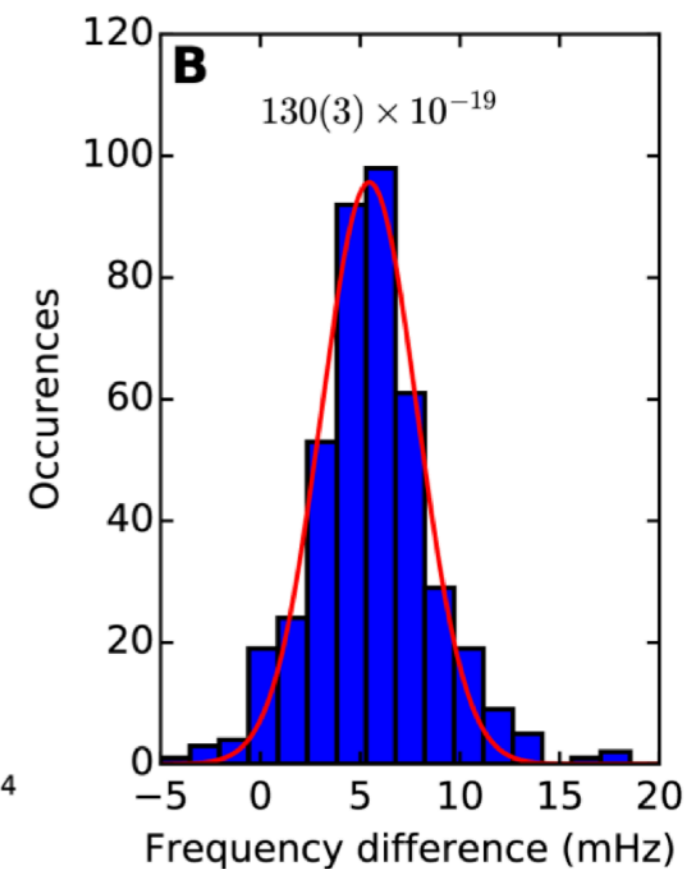
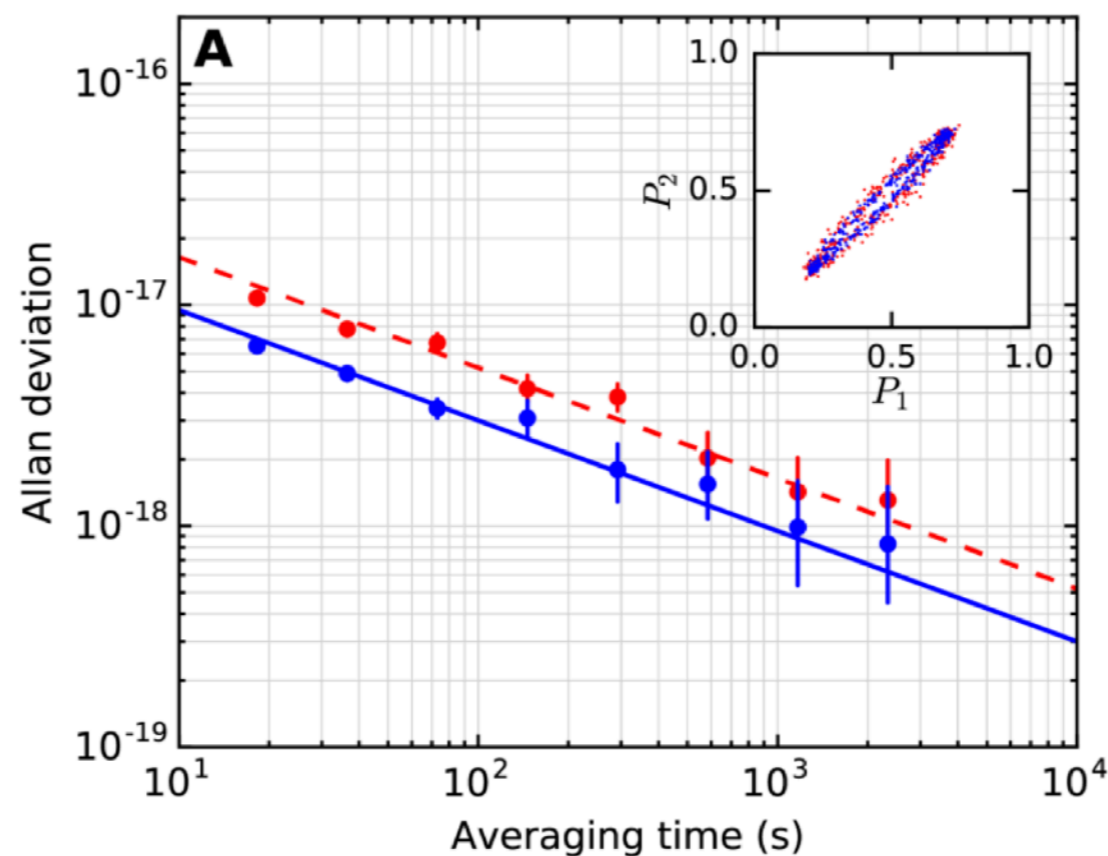


Tunnelling sensitive to the difference between Wannier and ground-state of harmonic trap, even in the tight-binding regime

# Deep lattices: Dicke regime and optical clocks



Precision  $< 10^{-18}$  over 1 hour  
 Campbell et al. *Science* **358**, 90 (2017)

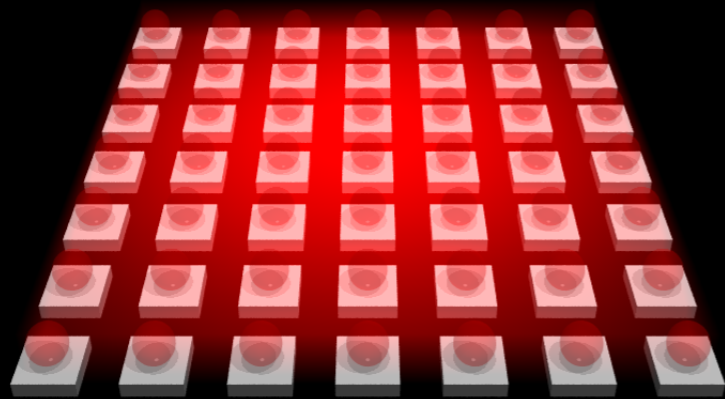


# Superfluid-to-Mott transition

## SUPERFLUID

$$J \gg U$$

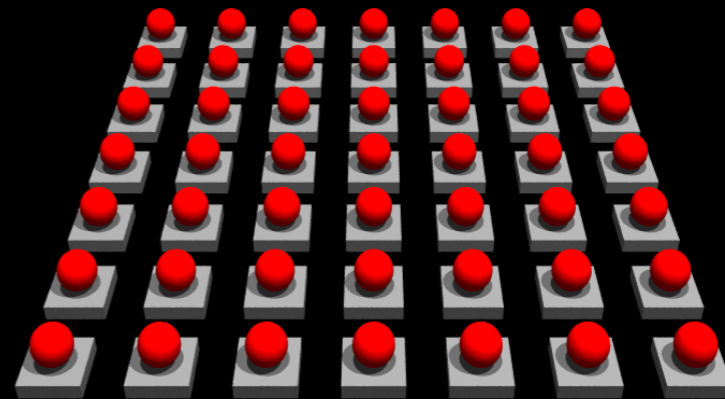
- ✓ Long-range phase coherence
- ✓ Poissonian number fluctuations
- ✓ Gapless excitation spectrum
- ✓ Compressible



## MOTT INSULATOR

$$U \gg J$$

- ✓ No phase coherence
- ✓ No number fluctuations (Fock states)
- ✓ Gap in the excitation spectrum
- ✓ Not compressible



Quantum phase transition (T=0) induced by interactions: **Mott transition**



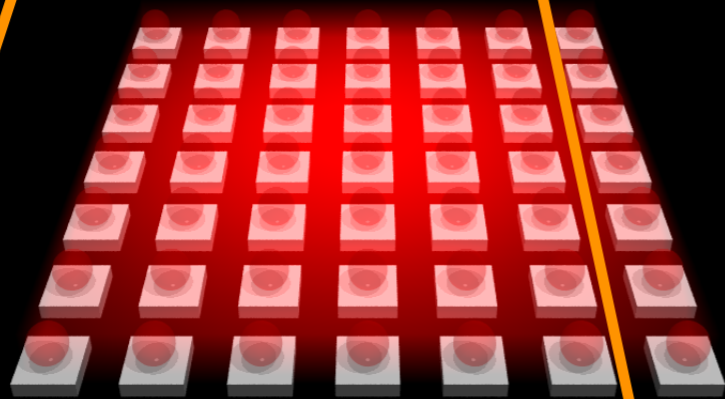
Sir N. Mott

# Superfluid-to-Mott transition

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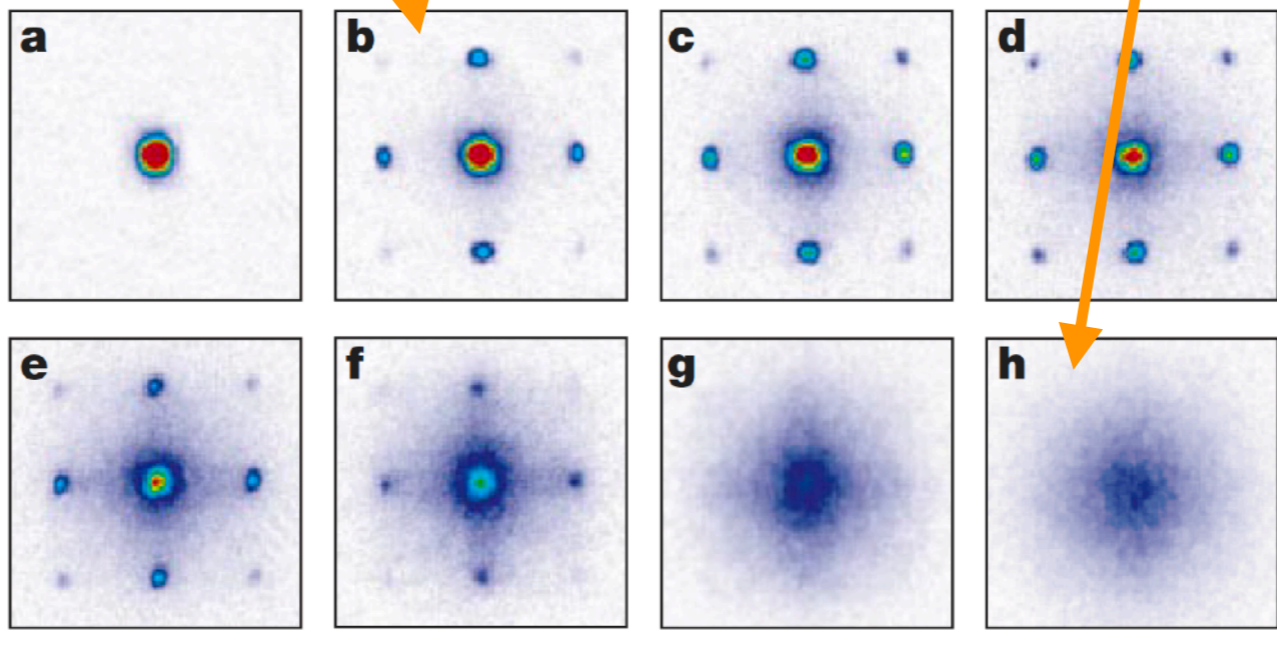
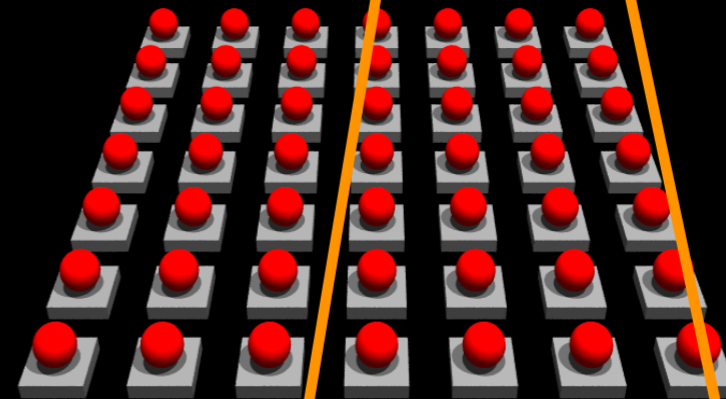
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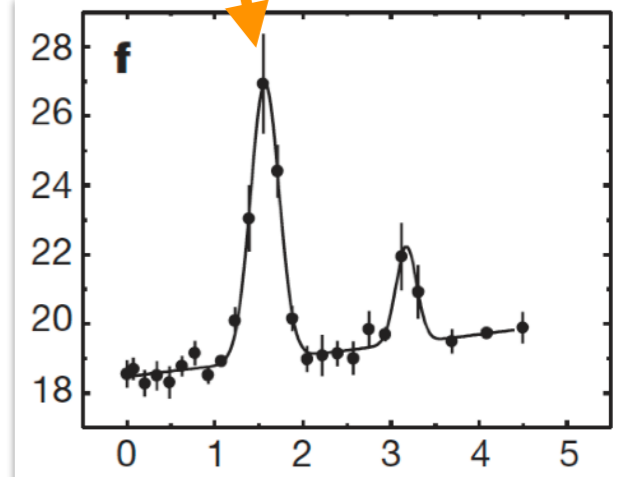
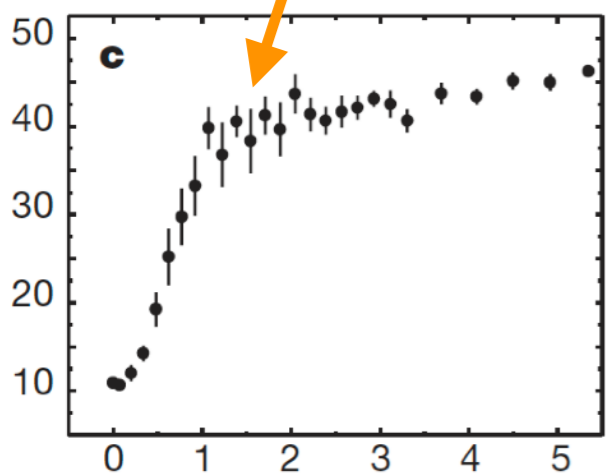
## MOTT INSULATOR

$$U \gg J$$

- ✓ No phase coherence
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- ✓ Not compressible



reversible loss of coherence (not shown)





# Locating the Mott transition

*(theoretical approaches)*

$\bar{n} = 1$	$\left(\frac{U}{t}\right)_c$ Mean-Field	$\left(\frac{U}{t}\right)_c$ Exact	
1D	11.66	3.37	DMRG
2D	23.33	16.74	QMC
3D	34.98	29.36	QMC

## Numerical techniques

- DMRG: density matrix renormalization group
- QMC: Quantum Monte Carlo

mean-field approaches over-estimates the extension of the order phase

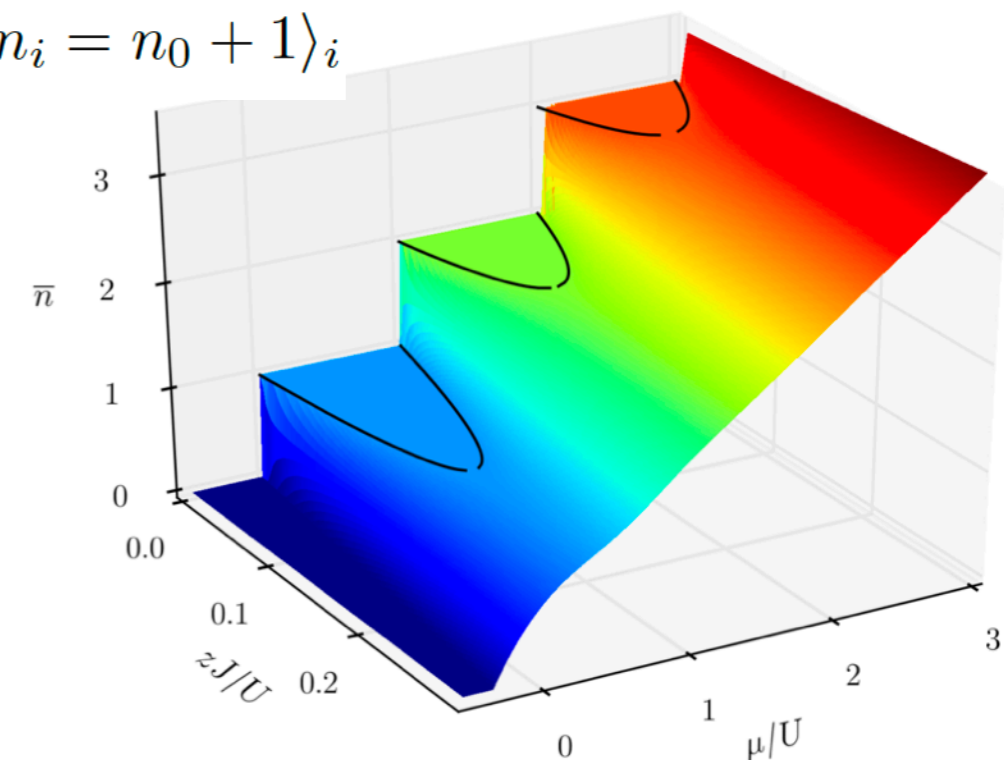
in low dimensional systems, quantum fluctuations play a larger role

**Alternative mean-field:** truncate Hilbert space to three states/site:

$$|\phi_i\rangle = c(n_0 - 1)|n_i = n_0 - 1\rangle_i + c(n_0)|n_i = n_0\rangle_i + c(n_0 + 1)|n_i = n_0 + 1\rangle_i$$

$$\mu_{n_0}^{(\pm)} = U \left( n_0 - \frac{1}{2} \right) \pm \frac{zJ}{2} \pm \sqrt{U^2 - 2UzJ(2n_0 + 1) + (zJ)^2}$$

<http://www.lkb.upmc.fr/boseeinsteincondensates/gerbier/>



# Locating the Mott transition

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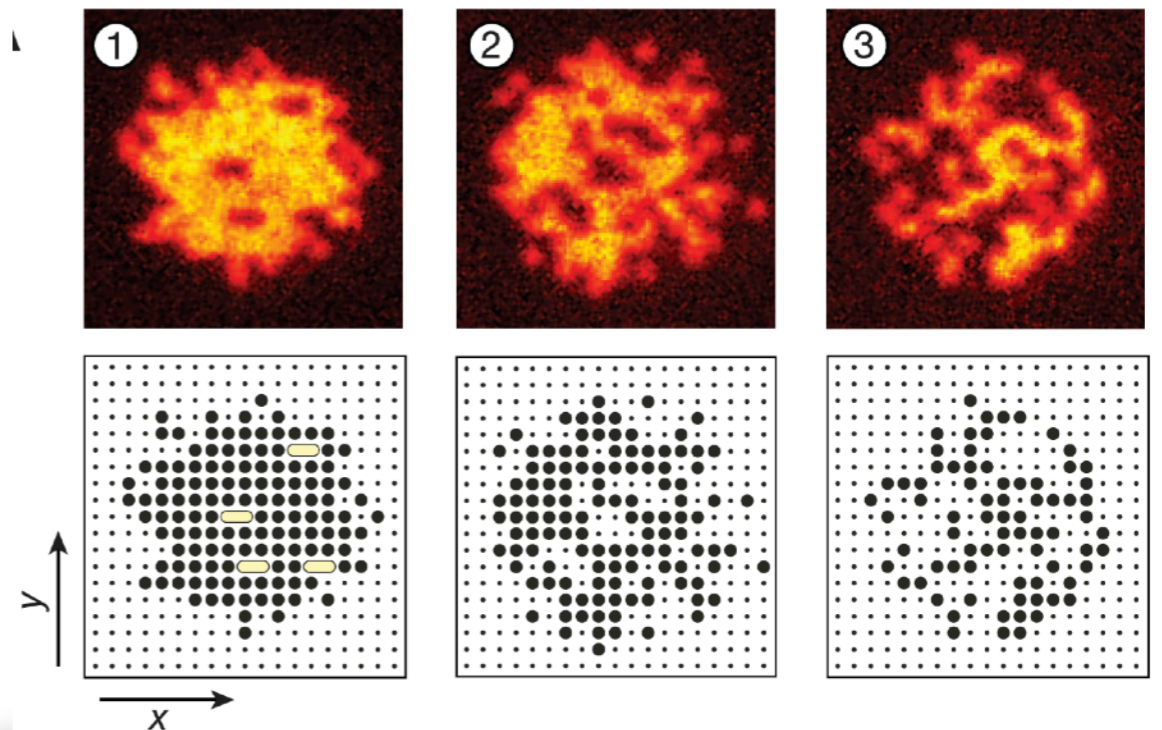
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Approach also "naturally" suited to investigate properties of the Mott state!  
*(quasi-particles: particle-hole excitations)*

$$J/U = 0.06$$

$$J/U = 0.11$$

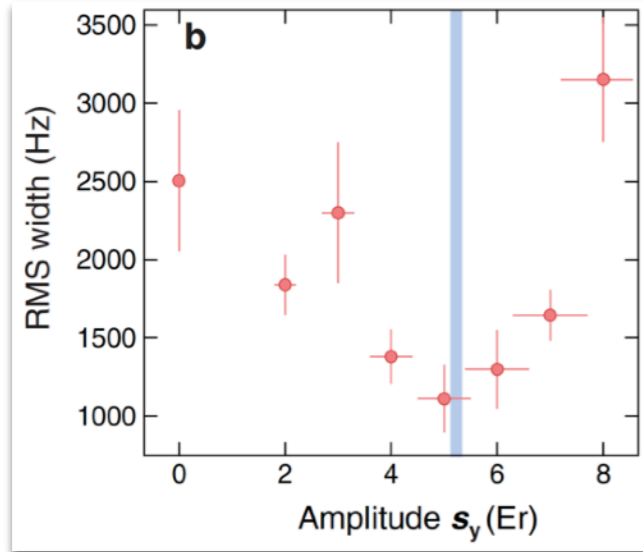
$$J/U = 0.3$$



# Locating the Mott transition

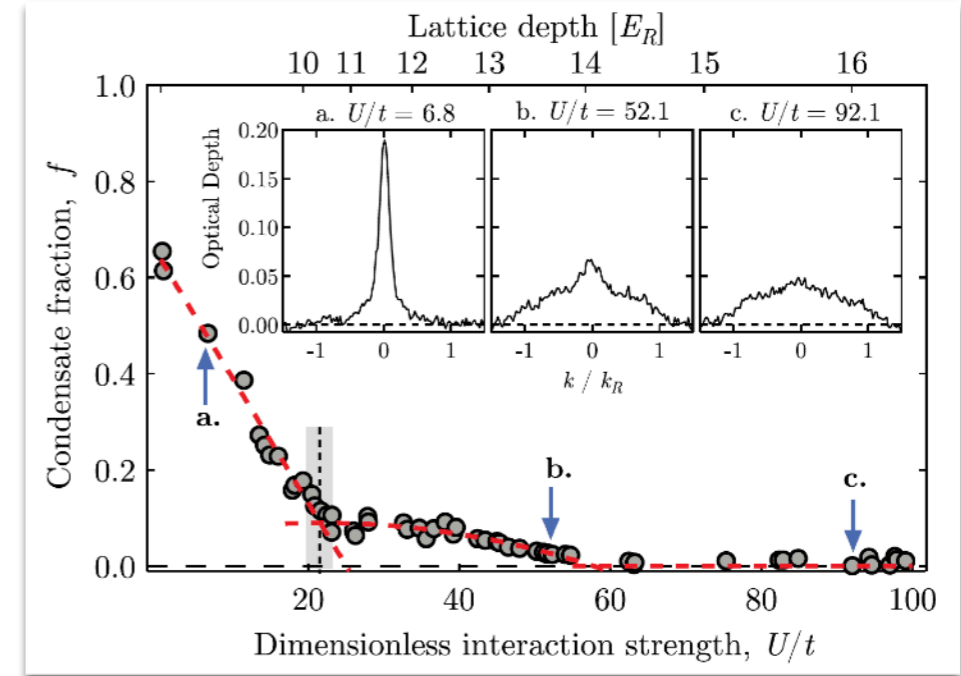
(experimental approaches)

1D - Bragg spectroscopy compatible with QMC prediction (not accurate though)

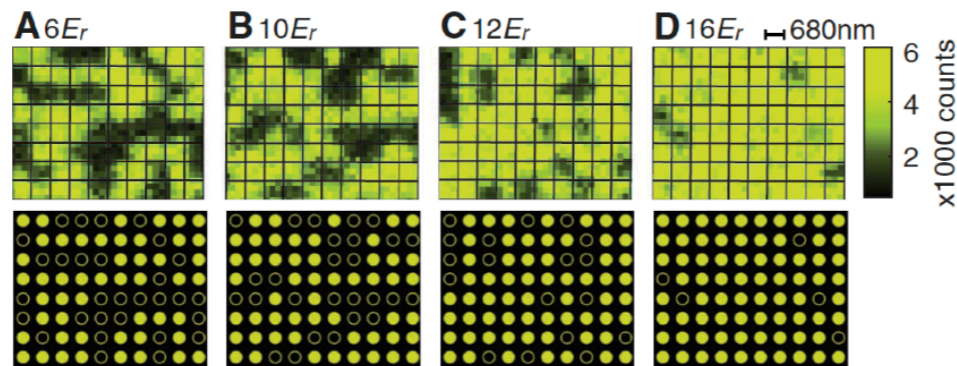


Clément et al. *Phys. Rev. Lett.* **102**, 155301 (2009)

2D - BEC fraction compatible with QMC prediction (not accurate though)

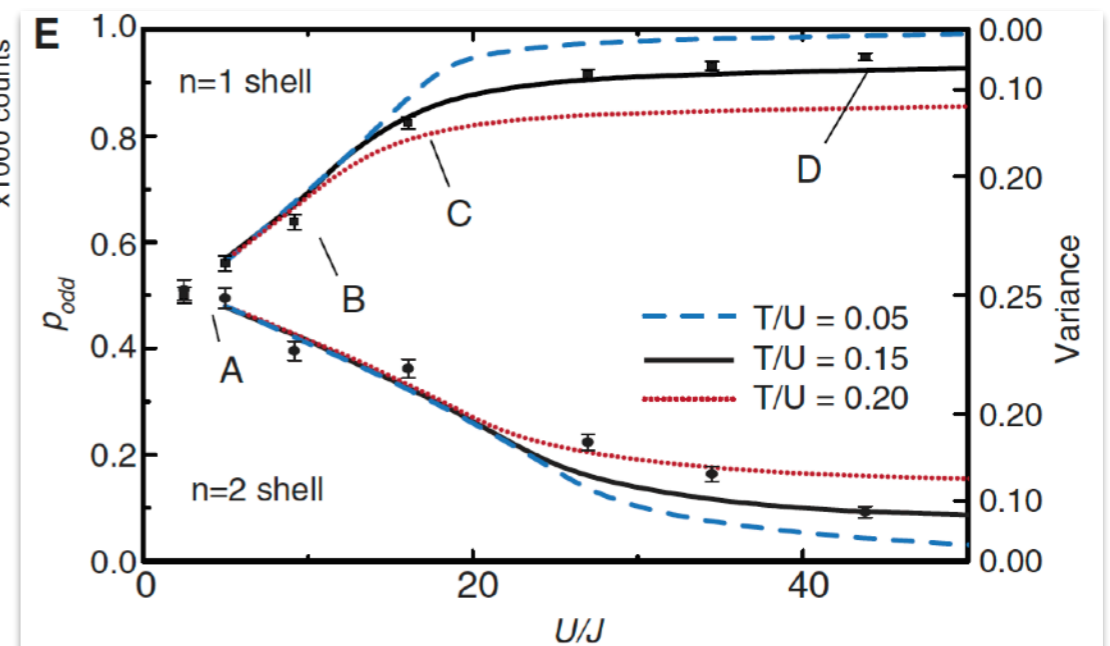


Jimenez-Garcia et al. *Phys. Rev. Lett.* **105**, 110401 (2010)



2D - Reduced atom number fluctuations in the Mott regime (not accurate though)

(parity measured)

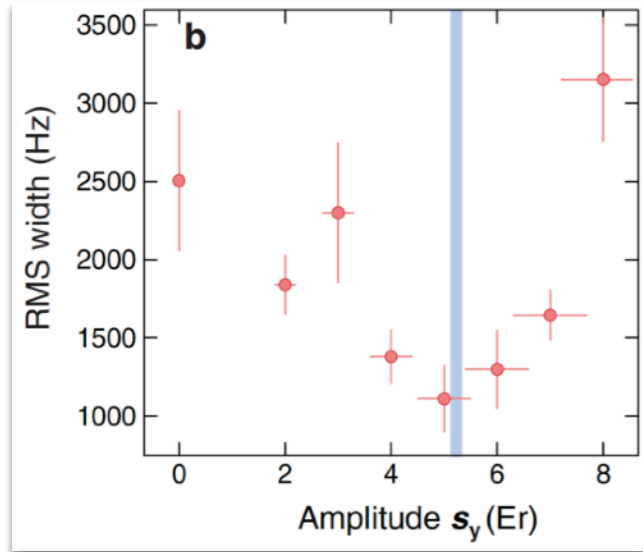


Bakr et al. *Science* **329**, 547 (2010)

# Locating the Mott transition

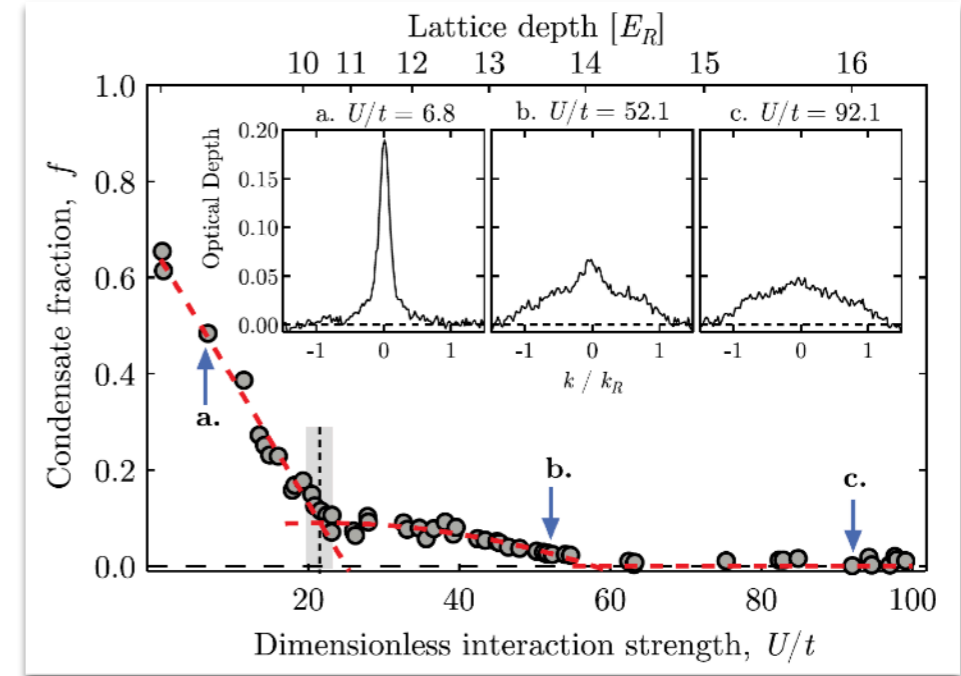
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Clément et al. *Phys. Rev. Lett.* **102**, 155301 (2009)

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Jimenez-Garcia et al. *Phys. Rev. Lett.* **105**, 110401 (2010)

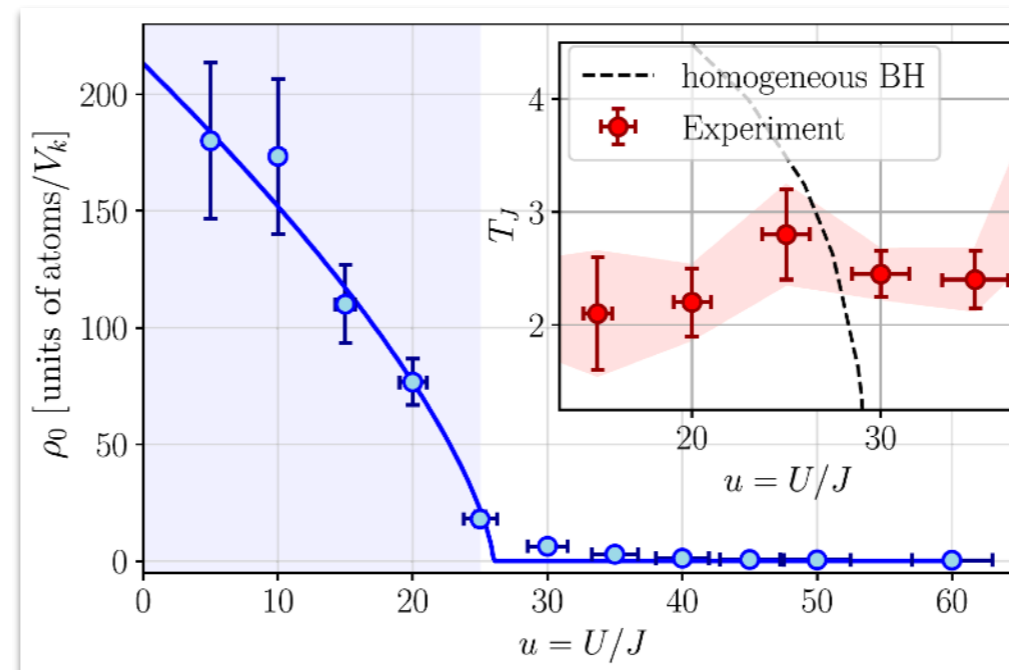
3D - different measures compatible with mean-field predictions...

Mun et al. *Phys. Rev. Lett.* **99**, 150604 (2007)

Becker et al. *New J. Phys.* **12**, 065025 (2010)

Mark et al. *Phys. Rev. Lett.* **107**, 175301 (2011)

Thomas et al. *Phys. Rev. Lett.* **119**, 100402 (2017)



Hercé et al. *Phys. Rev. A* **104**, L011301 (2021)

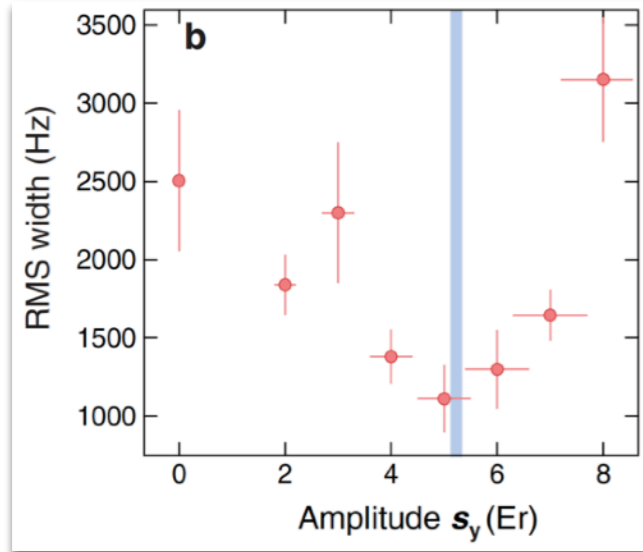
compatible with QMC (exclude mean-field)!

...but....

# Locating the Mott transition

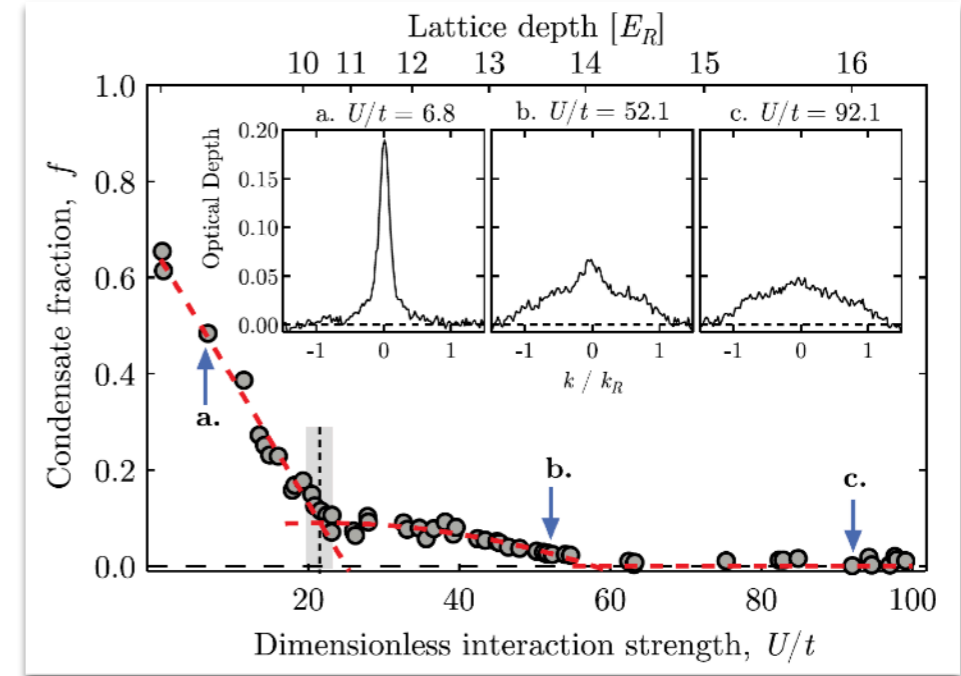
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*Phys. Rev. Lett.* **105**, 110401 (2010)

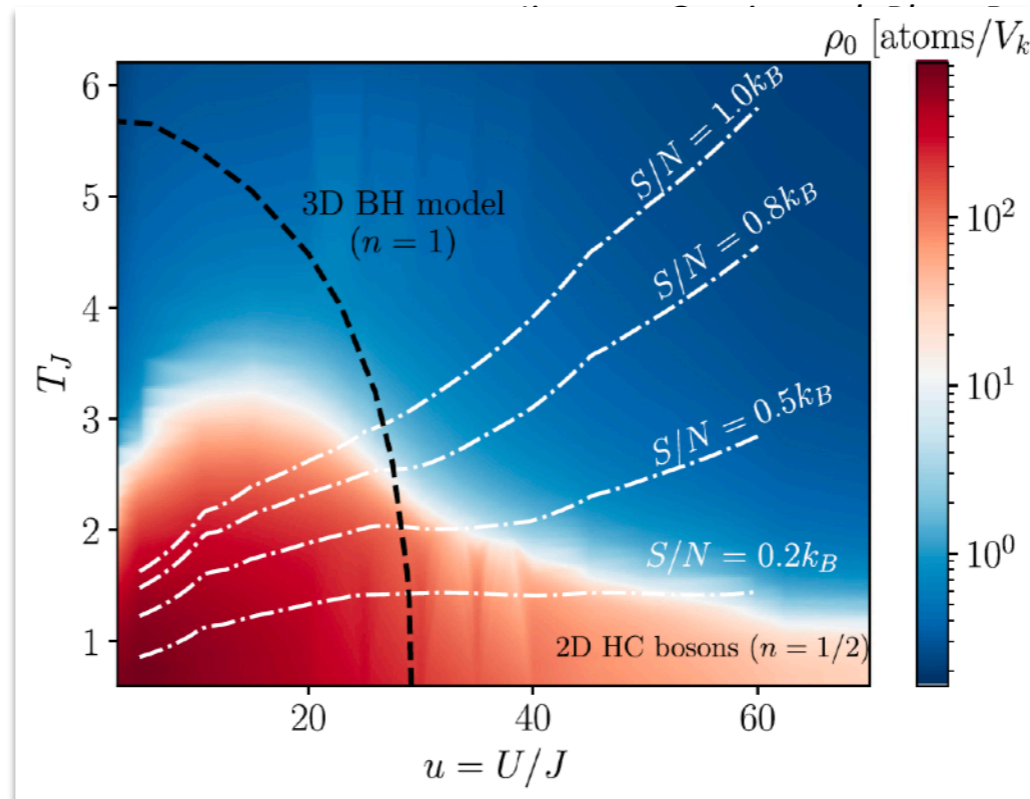
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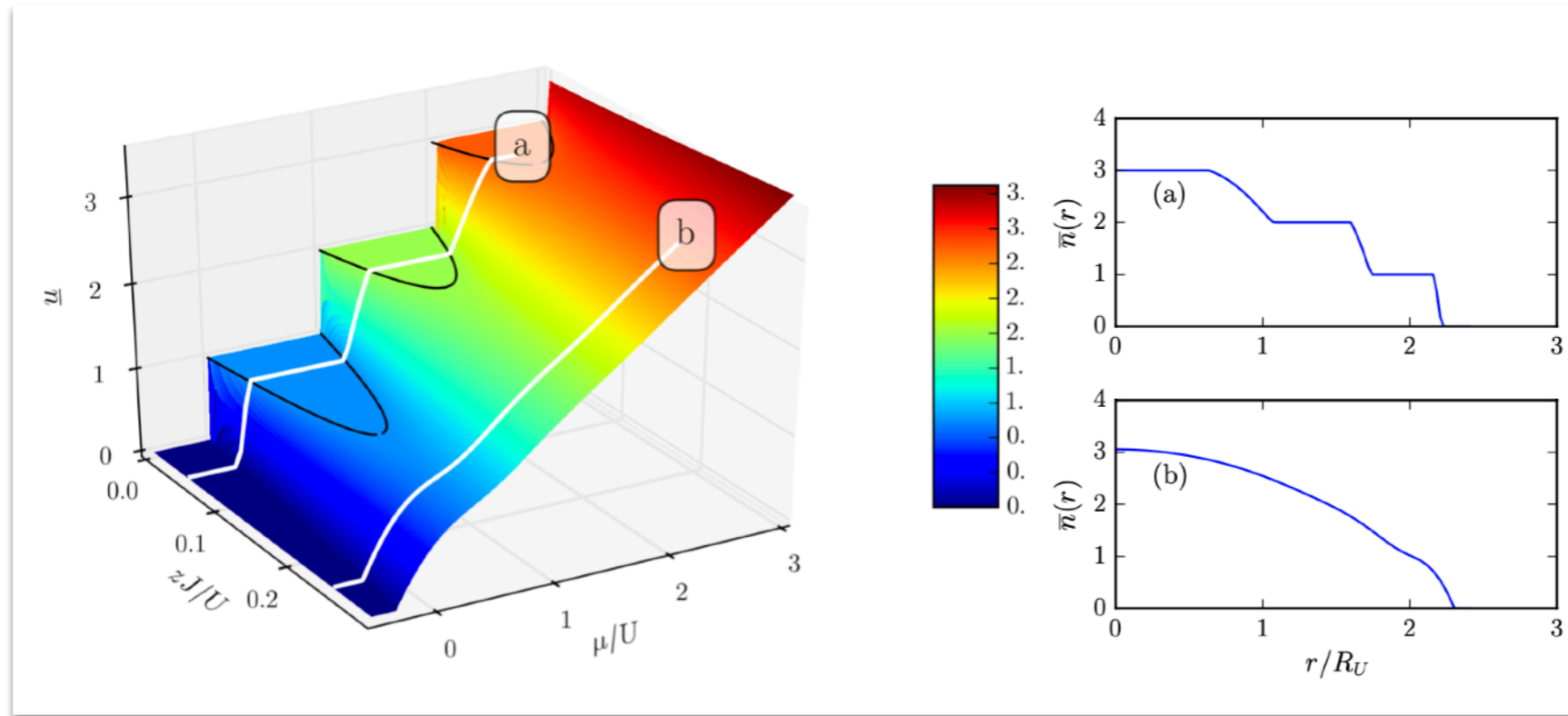
compatible with QMC (exclude mean-field)!

...but....

strongly depends on exp. parameters

# Inhomogeneous lattice Bose gases

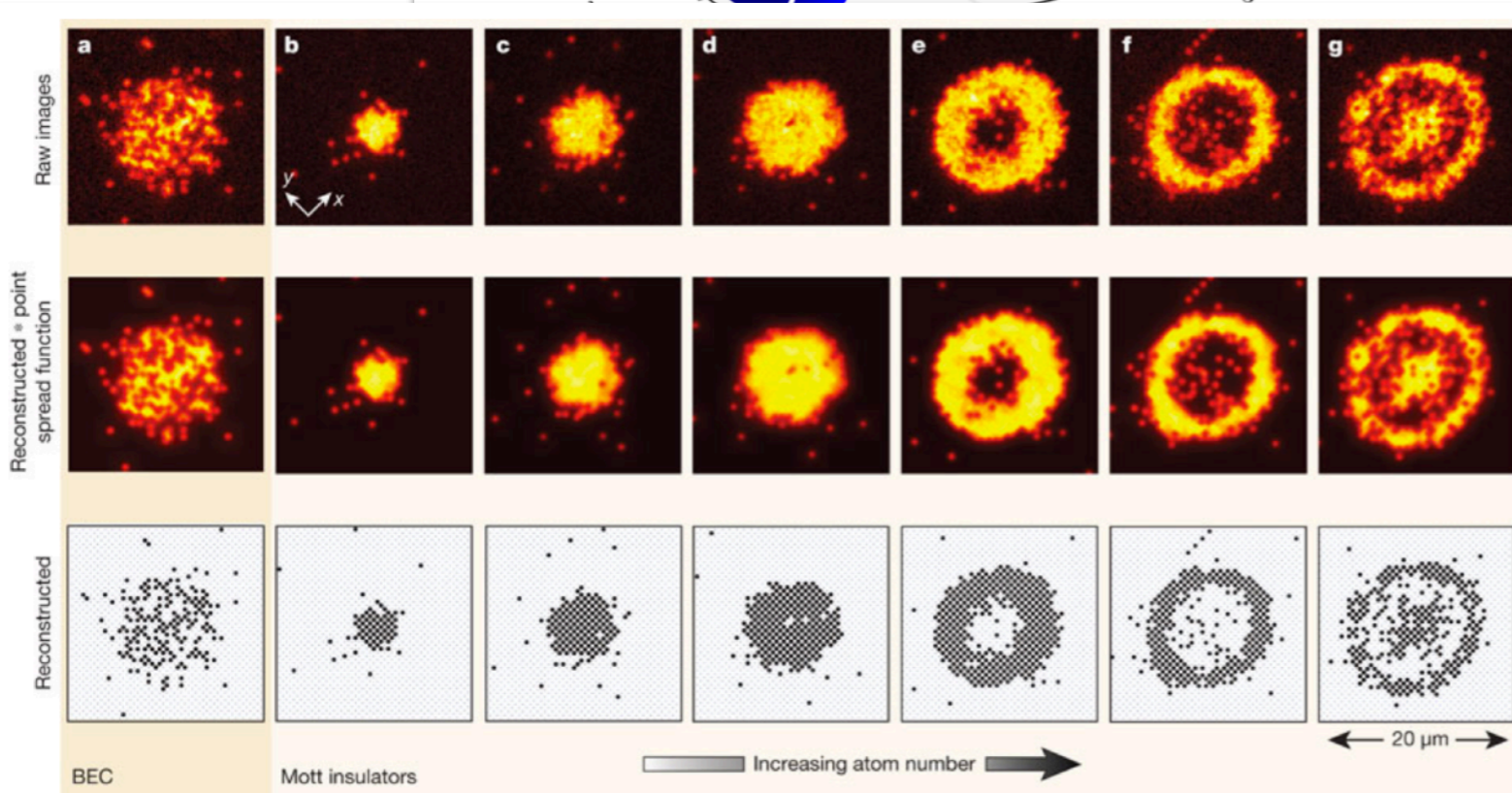
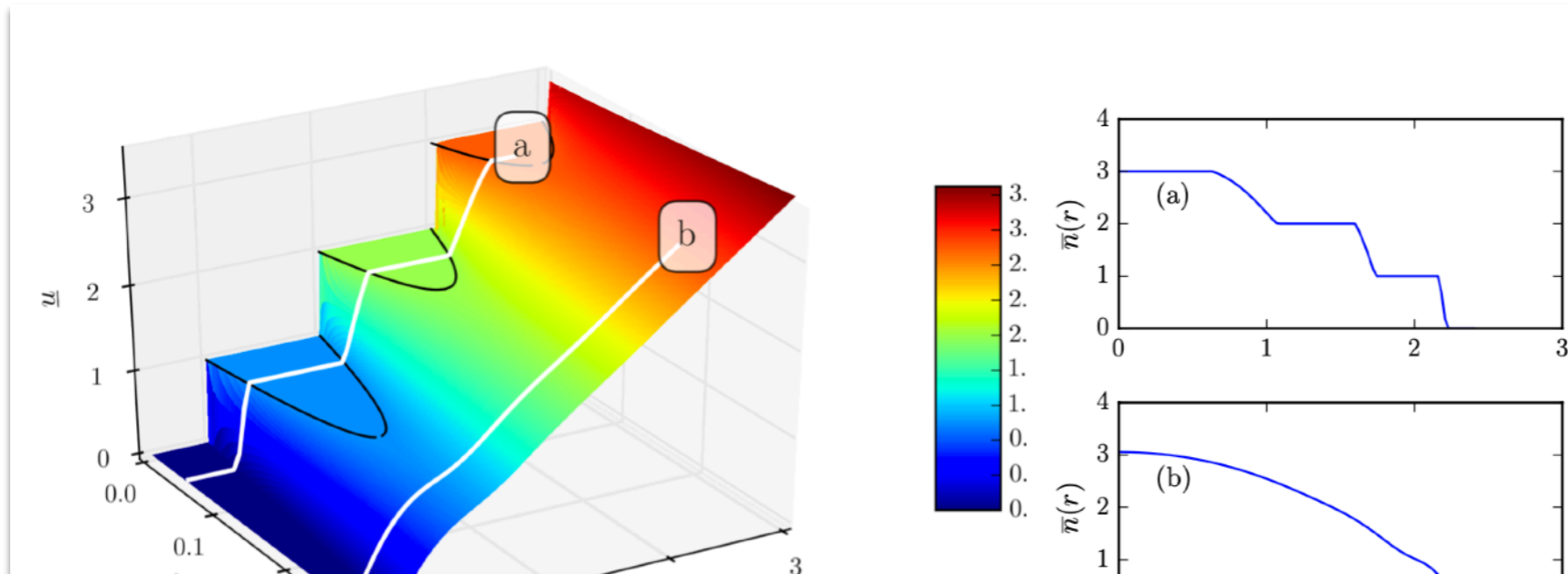
$$H = -J \sum_{j,j'} b_j^\dagger b_{j'} + \frac{U}{2} \sum_j n_j(n_j - 1) + \sum_j \frac{m\omega a^2}{2} j^2$$



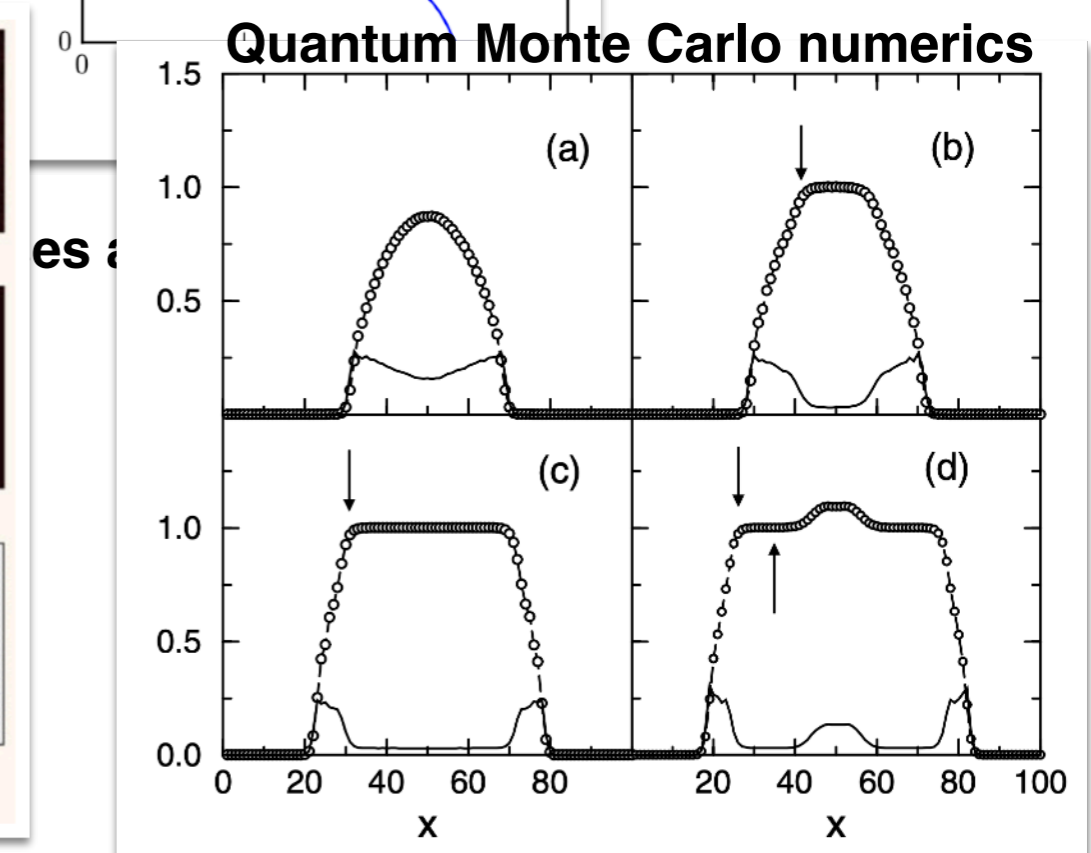
**co-existence of several phases at once!**

# Inhomogeneous lattice Bose gases

$$H = -J \sum_{j,j'} b_j^\dagger b_{j'} + \frac{U}{2} \sum_j n_j(n_j - 1) + \sum_j \frac{m\omega a^2}{2} j^2$$

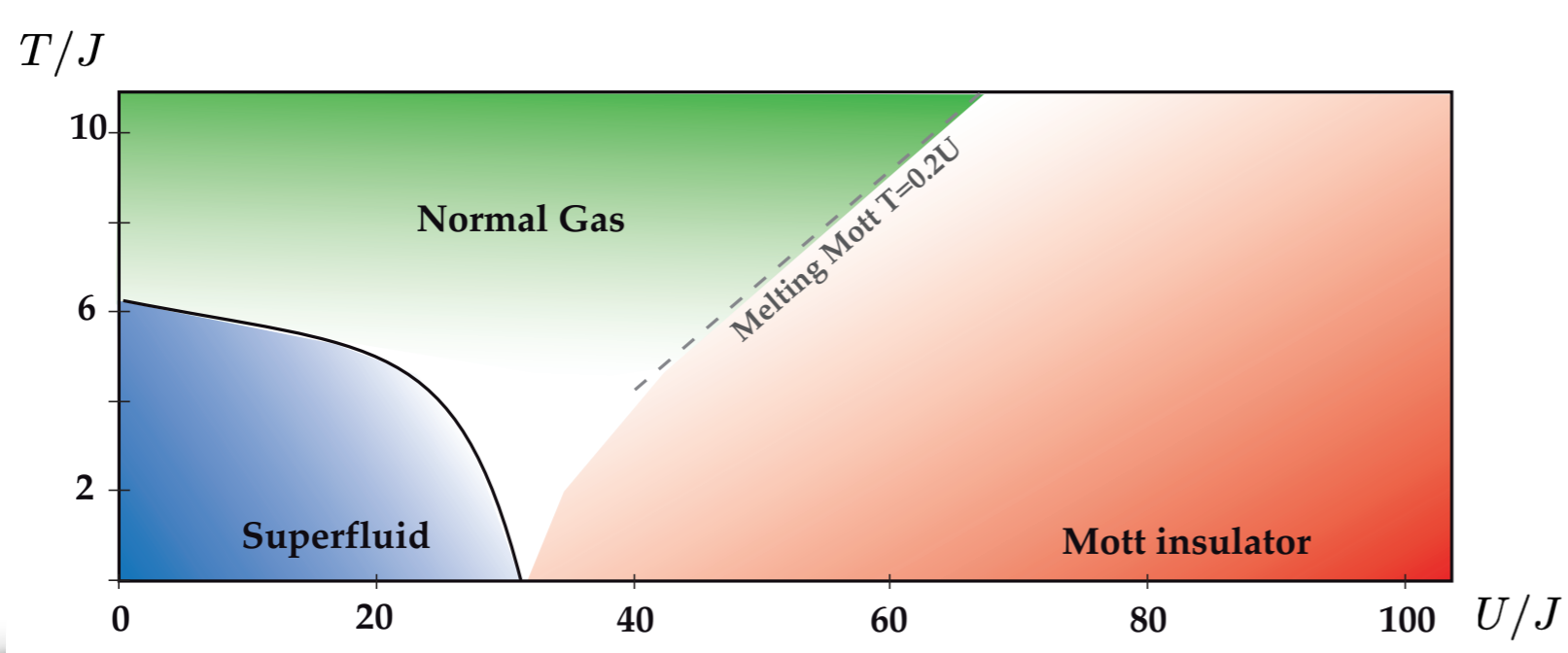


Sherson, *et al.* Nature **467**, 68-72 (2010)



G. Batrouni, *et al.* Phys. Rev. Lett. **89**, 117203 (2002)

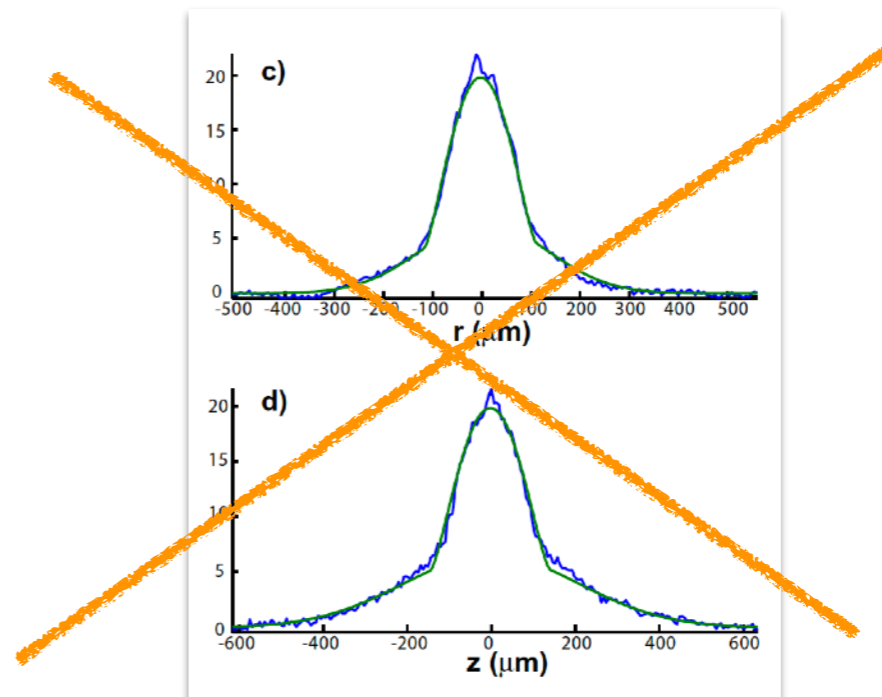
# Thermometry of lattice Bose gases



Experiments are conducted at finite temperature!

→ **observation of a finite temperature crossover rather than the quantum phase transition ( $T=0$ )**

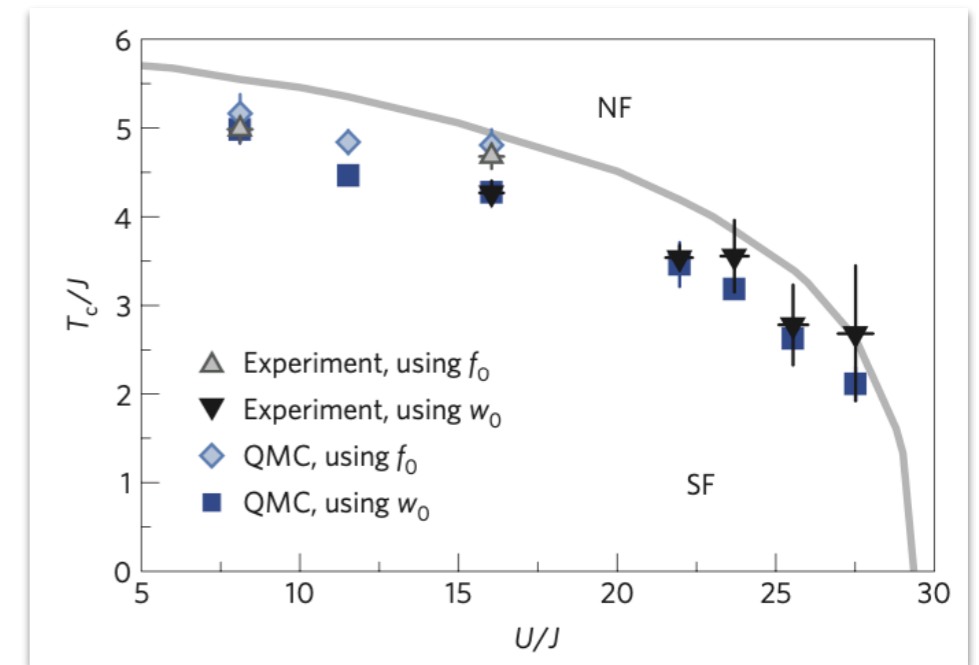
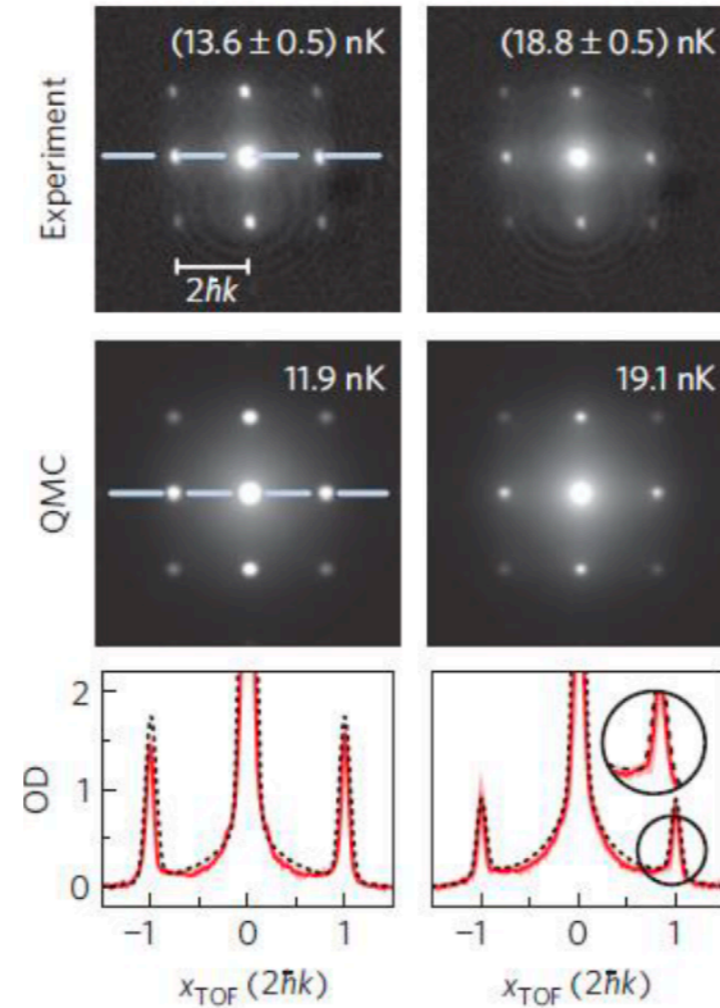
How to measure the temperature?





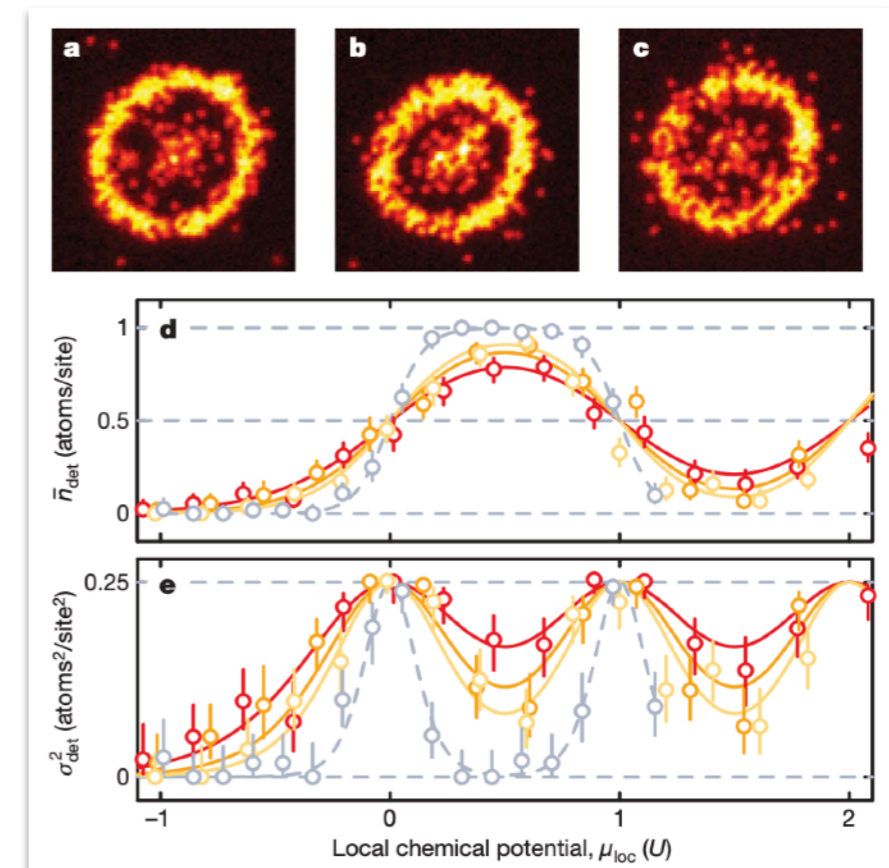
# Thermometry of lattice Bose gases

3D - comparison with ab-initio QMC calculations after TOF



Trotzky et al. *Nat. Phys.* **6**, 998 (2010)

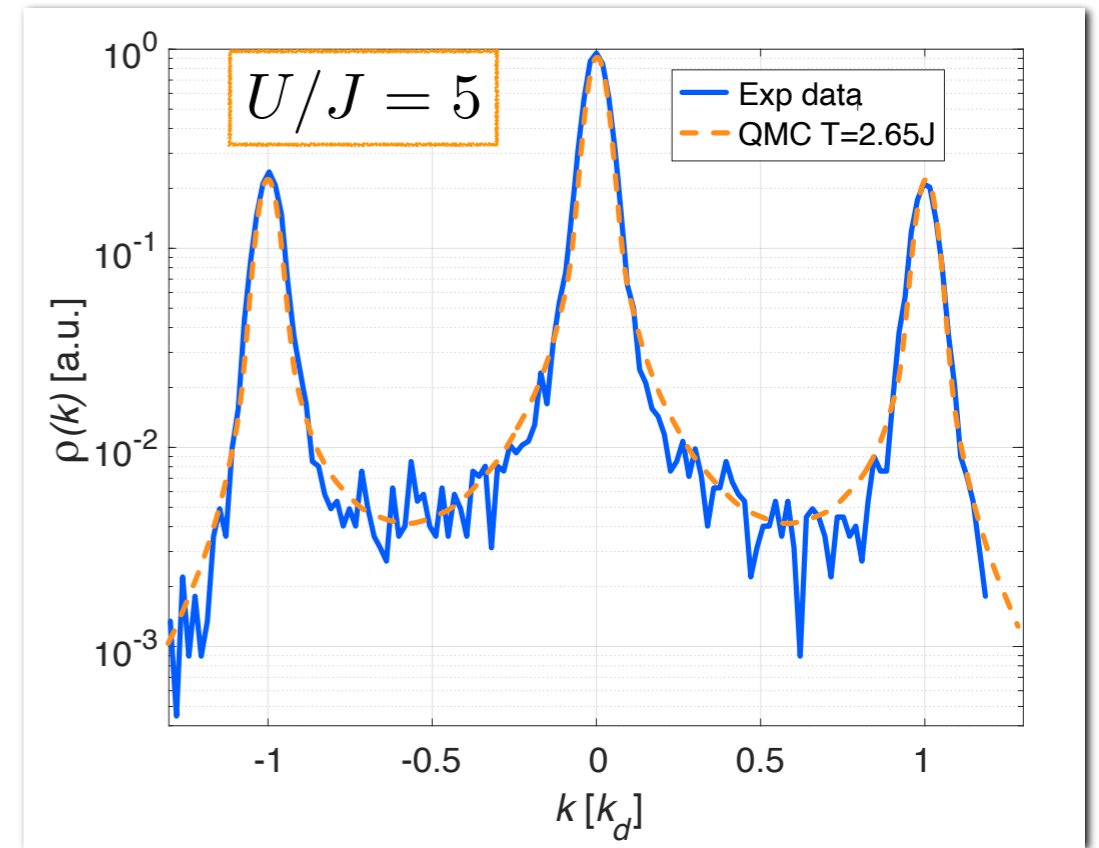
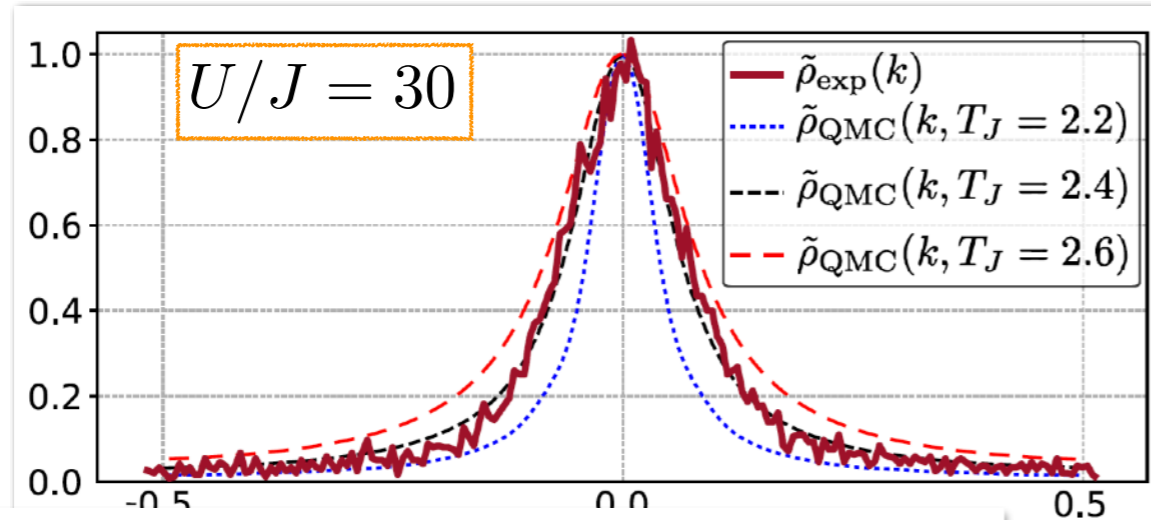
2D - site-resolved atom number fluctuations in the Mott regime



Sherson et al. *Nature* **467**, 68 (2010)

# Thermometry of lattice Bose gases

3D - comparison with ab-initio QMC calculations after TOF

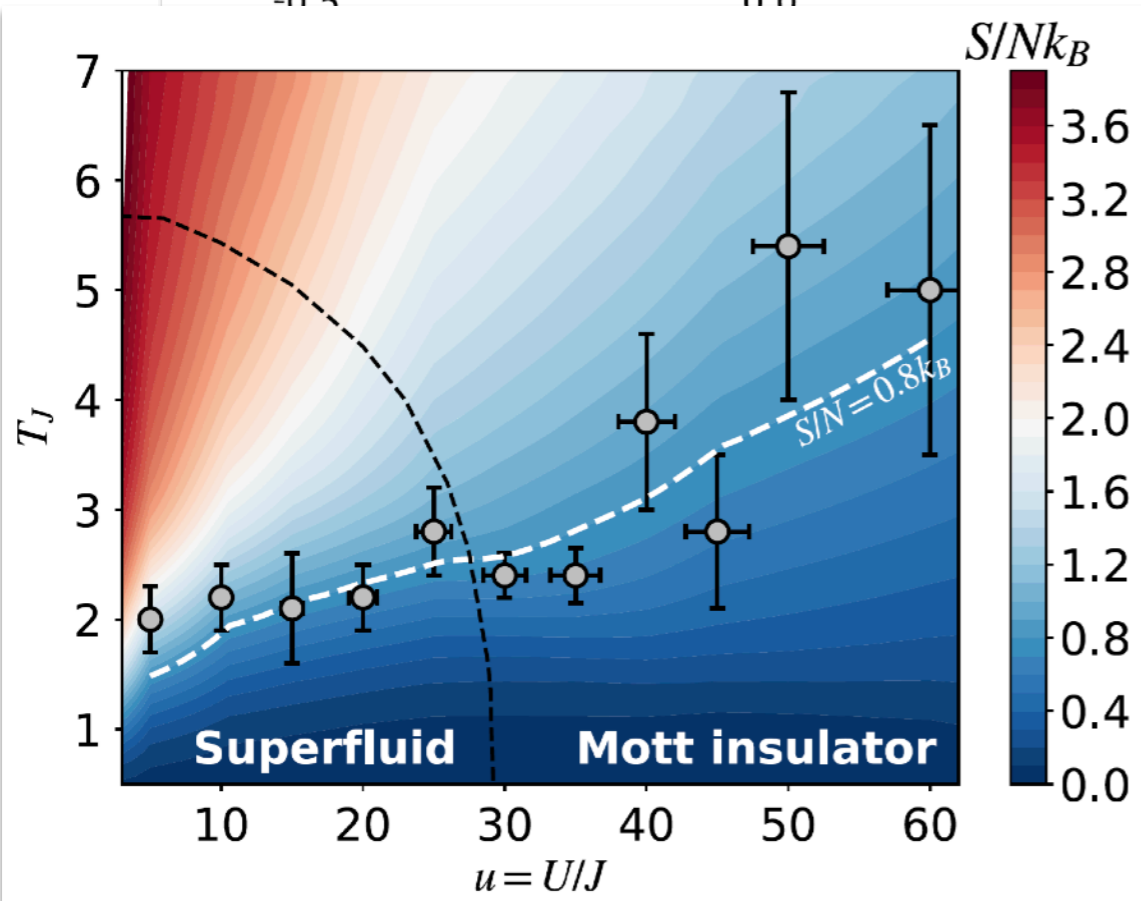


Cayla et al. Phys. Rev. A **97**, 061609 (2018)

3D - comparison with ab-initio QMC calculations after long TOF in the entire phase diagram: **compatible with adiabatic loading**

saturates Cramer-Rao bound everywhere!

$$\delta T = \frac{1}{\sqrt{I(T)M}}$$



Carcy et al. Phys. Rev. Lett. **126**, 045301 (2021)